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Corrigendum to "Asymptotics for the Unconstrained Polarization (Chebyshev) Problem"

Douglas Hardin¹ · Mircea Petrache² · Edward B. Saff¹ ©

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The purpose of this note is to address two needed modifications in the recently published article [3]; namely, revisions to Theorems 1.11 and 1.12 and their proofs.

In the revised Theorem 1.11, we add the condition in case (i) that A be strongly d-rectifiable and, in case (ii) we now require s > p-2 and replace the condition on the regularity of the energy equilibrium measure μ_s^A by the requirement that $\operatorname{supp}(\mu_{s,A}) = A$ and that the potential $\mathcal{U}_s(\mu_{s,A},x)$ is constant on A, where for a positive measure ν on \mathbb{R}^p and s > 0:

$$\mathcal{U}_s(v,x) := \int \frac{1}{|x-y|^s} dv(y). \tag{1}$$

When ω is a finite set in \mathbb{R}^p , we also identify ω with its counting measure and write

$$\mathcal{U}_s(\omega, x) := \sum_{y \in \omega} \frac{1}{|x - y|^s}.$$

Theorem 1 (Replacement of [3, Thm. 1.11]) For integers p, d such that $p \ge 2$, $1 \le d \le p$ and $A \subset \mathbb{R}^p$ compact, suppose that one of the following conditions holds:

(i) $s > \max\{d, p-2\}$ and A is strongly d-rectifiable with $\mathcal{H}_d(A) > 0$.

☐ Edward B. Saff Edward.B.Saff@Vanderbilt.edu

Douglas Hardin douglas.hardin@Vanderbilt.edu

Mircea Petrache decostruttivismo@gmail.com

- Vanderbilt University, Nashville, TN, USA
- Pontificia Universidad Catolica de Chile, Santiago, Chile



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(ii) p-2 < s < d, the set A is d-regular and the equilibrium measure $\mu_{s,A}$ satisfies $\operatorname{supp}(\mu_{s,A}) = A$ and $\mathcal{U}_s(\mu_{s,A},x)$ is constant for $x \in A$. Then both limits $h_{s,d}(A)$ and $h_{s,d}^*(A)$ exist, are finite, and

$$h_{s,d}(A) = h_{s,d}^*(A).$$
 (2)

The revision of Theorem 1.12 requires in the case s = p that A be Jordan measurable, but no change for the s > p case.

Theorem 2 (Replacement of [3, Thm. 1.12]) If $A \subset \mathbb{R}^p$ is a compact set and s > p, or if s = p and $\mathcal{L}^p(\partial A) = 0$, then

$$h_{s,p}^*(A) = h_{s,p}(A) = \frac{\sigma_{s,p}}{\mathcal{L}_p(A)^{s/p}}.$$
 (3)

Moreover, if $\mathcal{L}_p(A) > 0$, then for any asymptotically extremal sequence $\Omega = \{\omega_N\}_{N \geq 1}$ (for either the constrained or unconstrained polarization problem) we have the weak-* convergence

$$\frac{1}{N} \sum_{x_i \in \omega_N} \delta_{x_i} \stackrel{*}{\rightharpoonup} \frac{\mathcal{L}_p \lfloor A}{\mathcal{L}_p(A)} \quad as \quad N \to \infty, \tag{4}$$

where $\mathcal{L}_p[A := \mathcal{L}_p(\cdot \cap A)$ is the restriction to A of \mathcal{L}_p .

We first address Theorem 1 and its proof. We thank Alexander Reznikov who brought to our attention that the derivation of the inequality (4.20) in the proof of [3, Thm. 1.11] requires modifications in the case p > d. Under the assumptions of Theorem 1, we now conclude that the limit $h_{s,d}^*(A)$ exists and is finite (in contrast, this was only established in part (ii) of [3, Thm. 1.11]).

Proof of Theorem 1 We first discuss the subcase p = d of case (i). We note that the proof of [3, Thm. 1.11] remains valid for this subcase; that is, if $\mathcal{L}_p(A) > 0$, s > p = d and $h_{s,p}^*(A)$ exists, then $h_{s,p}(A)$ also exists and equals $h_{s,p}^*(A)$. In the corrected proof given below for Theorem 2 we establish that in this subcase $h_{s,p}^*(A)$ exists and is equal to the rightmost expression in (3). In particular it is finite if $\mathcal{L}_p(A) > 0$. Hence (2) holds for this subcase.

To complete the proof of case (i) for p > d, we remark that the proof of [3, Thm. 1.14] relies on the s > p case of [3, Thm. 1.12], which coincides with the s > p case of Theorem 2. Thus the proof of [3, Thm. 1.14] remains valid. Then (2) follows from [3, Thm. 1.14].

It remains to consider the case (ii). We will use the following notation. For 0 < s < p, $\omega \subset \mathbb{R}^p$ finite and ν a positive finite measure over \mathbb{R}^p , let $P_s(A, \nu) := \min_{v \in A} \mathcal{U}_s(\nu, y)$.

Our first observation is that under the assumptions that 0 < s < d and A d-regular, the proof of [5, Thm. 1.8] extends under weaker hypotheses as stated in the following result.

Theorem 3 (Extension of [5, Thm. 1.8]) Let f be a d-Riesz-like function and $A \subset \mathbb{R}^p$ be a d-regular compact set. If ω_N is a sequence of N-point configurations in \mathbb{R}^p such that

$$\nu_N := \frac{1}{N} \sum_{x \in \omega_N} \delta_x \stackrel{*}{\rightharpoonup} \nu$$

for a probability measure ν supported on A, then $P_f(A,\nu_N) \to P_f(A,\nu)$.

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For the definitions of *d-Riesz-like potentials* f (which include Riesz potentials with power 0 < s < d) and the f-polarization $P_f(A, \mu)$ of a measure μ over the set A, see [5, Defs. 1.1 and 1.7]. Note that $P_s(A, \mu) = P_f(A, \mu)$ with $f(|x|) = |x|^{-s}$ for s > 0.

Proof of Theorem 3 The original version of [5, Thm. 1.8] required that $\omega_N \subset A$, however this hypothesis is not used since the main potential theoretic result needed in the proof is the principle of descent [5, Thm. 2.3], which works in general. On the other hand, the fact that ν is supported on A is required in order to use the minimum principle for Riesz-like f, see [5, Thm. 2.5].

Proof of Theorem 1 (ii) continued. Suppose now that the hypotheses of (ii) hold. For $N \ge 1$, let ω_N^A be such that

$$P_s(A, \omega_N^A) = \mathcal{P}_s^*(A, N). \tag{5}$$

By [3, Cor. 1.9] if the empirical measures of a subsequence ω_N^A , $N \in \mathcal{N}$, as in (5) converge weakly to a measure ν , then ν is supported on A. By Theorem 3, it follows that

$$\lim_{N \in \mathcal{N}} \frac{\mathcal{P}_s^*(A, N)}{N} = P_s(A, \nu). \tag{6}$$

Furthermore, by [4, Thm. 2], ν must maximize the continuous (integral) unconstrained polarization, which therefore coincides with the maximal continuous constrained polarization (since ν is supported on A):

$$P_s(A, \nu) = \max_{\mu \in \mathcal{P}(\mathbb{R}^p)} P_s(A, \mu) = \max_{\mu \in \mathcal{P}(A)} P_s(A, \mu). \tag{7}$$

We can then apply the result [6, Thm. 1.2], valid for constrained polarization, which implies $v = \mu_{s,A}$. In particular, v has no atoms since it has finite s-energy.

Let $\operatorname{proj}_A : \mathbb{R}^p \to A$ be a measurable map which assigns to any point $x \in \mathbb{R}^p$ a point $y \in A$ such that $|x - y| = \operatorname{dist}(x, A)$. Then for fixed $\epsilon > 0$ we set

$$\omega_{N,\epsilon}^A := \mathrm{proj}_A(\omega_N^A \cap A_\epsilon), \qquad \nu_{N,\epsilon} := \frac{1}{N} \sum_{x \in \omega_{N,\epsilon}^A} \delta_x.$$

Let $y_{N,\epsilon} \in A$ be a point at which $\mathcal{U}_s(v_{N,\epsilon}, y)$ achieves its minimum over A. We then fix a sequence $\mathcal{E} = \{\epsilon_k\}$ with $\epsilon_k \to 0$ and an increasing sequence $\mathcal{N}^0 \subset \mathbb{N}$. Up to taking a subsequence, we may assume that

$$\nu_N := \frac{1}{N} \sum_{x \in \omega_N^A} \delta_x \stackrel{*}{\rightharpoonup} \nu.$$

By the discussion in the preceding paragraph we see that (6) and (7) hold for such ν , and we have $\nu = \mu_{s.A}$.

By the compactness of A, via a diagonalization procedure we can find an increasing subsequence $\mathcal{N} \subset \mathcal{N}^0$ such that for each $\epsilon \in \mathcal{E}$ there holds $\lim_{N \in \mathcal{N}} y_{N,\epsilon} = y_{\infty,\epsilon}$. Up to taking a subsequence of \mathcal{E} we can further assume that $\lim_{\epsilon \in \mathcal{E}} y_{\infty,\epsilon} = y_{\infty}$.

Now fix $\delta > 0$. If we let

$$\omega_{N,\epsilon,\delta}^A := \operatorname{proj}_A(\omega_N^A \cap A_\epsilon \setminus B(y_\infty, \delta)), \quad \nu_{N,\epsilon,\delta} := \frac{1}{N} \sum_{x \in \omega_{N,\epsilon,\delta}^A} \delta_x,$$

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then we can write

$$\mathcal{U}_{s}(\nu_{N,\epsilon,\delta}, y_{N,\epsilon}) - \mathcal{U}_{s}(\nu_{N} \lfloor (A_{\epsilon} \setminus B(y_{\infty}, \delta)), y_{N,\epsilon})$$

$$= \frac{1}{N} \sum_{x \in \omega_{N,\epsilon}^{A} \cap A_{\epsilon} \setminus B(y_{\infty}, \delta)} \left(\frac{1}{|y_{N,\epsilon} - \operatorname{proj}_{A}(x)|^{s}} - \frac{1}{|y_{N,\epsilon} - x|^{s}} \right).$$

Now note that for all but finitely many indices of the subsequences \mathcal{N} and \mathcal{E} there holds $|y_{N,\epsilon} - y_{\infty}| < \delta/3$ and for any $x \in \omega_N^A \cap A_{\epsilon} \setminus B(y_{\infty}, \delta)$ there holds $|\operatorname{proj}_A(x) - x| \le \epsilon < \delta/6$. By triangular inequality, for such N, ϵ, x we get

$$|y_{N,\epsilon} - \operatorname{proj}_{A}(x)| > \delta/2, \quad |y_{N,\epsilon} - x| > \delta/2, \quad \left| \frac{1}{|y_{N,\epsilon} - \operatorname{proj}_{A}(x)|^{s}} - \frac{1}{|y_{N,\epsilon} - x|^{s}} \right| < C \frac{\epsilon}{\delta^{s+1}},$$
(8)

where for the last above estimate we have used the fact that $|x - y|^{-s}$ is regular over the region $|x - y| > \delta/2$. From the above it follows that for each fixed $\delta > 0$ there holds

$$\lim_{\epsilon \in \mathcal{E}} \lim_{N \in \mathcal{N}} \left| \mathcal{U}_s(\nu_{N,\epsilon,\delta}, y_{N,\epsilon}) - \mathcal{U}_s(\nu_N \lfloor (A_\epsilon \setminus B(y_\infty, \delta)), y_{N,\epsilon}) \right| = 0.$$
 (9)

By using the positivity of our kernel and [3, Cor. 1.9], for all $\epsilon > 0$ there holds

$$0 \leq \lim_{N \to \infty} \mathcal{U}_s(\nu_N \lfloor (\mathbb{R}^p \setminus (A_\epsilon \cup B(y_\infty, \delta))), y_{N, \epsilon})$$

$$\leq \lim_{N \to \infty} \mathcal{U}_s(\nu_N \lfloor (\mathbb{R}^p \setminus A_\epsilon), y_{N, \epsilon}) \leq \epsilon^{-s} \lim_{N \to \infty} \frac{\#(\omega_N^A \setminus A_\epsilon)}{N} = 0.$$

It follows that all the above terms are zero. Observing that $\mathbb{R}^p \setminus B(y_\infty, \delta) = (\mathbb{R}^p \setminus (A_\epsilon \cup B(y_\infty, \delta))) \cup (A_\epsilon \setminus B(y_\infty, \delta))$, it follows by superposition that

$$\lim_{\epsilon \in \mathcal{E}} \lim_{N \in \mathcal{N}} \mathcal{U}_{s}(\nu_{N} \lfloor (A_{\epsilon} \setminus B(y_{\infty}, \delta)), y_{N, \epsilon}) = \lim_{\epsilon \in \mathcal{E}} \lim_{N \in \mathcal{N}} \mathcal{U}_{s}(\nu_{N} \lfloor (\mathbb{R}^{p} \setminus B(y_{\infty}, \delta)), y_{N, \epsilon}).$$
(10)

By the principle of descent, applied to the measures $\nu_N \lfloor (\mathbb{R}^p \setminus B(y_\infty, \delta))$ at the points $y_{N,\epsilon}$, and then to $\nu \lfloor (\mathbb{R}^p \setminus B(y_\infty, \delta))$ and points $y_{\infty,\epsilon}$, we obtain

$$\liminf_{\epsilon \in \mathcal{E}} \inf_{N \in \mathcal{N}} \mathcal{U}_{s}(\nu_{N} \lfloor (\mathbb{R}^{p} \setminus B(y_{\infty}, \delta)), y_{N, \epsilon}) \\
\geq \liminf_{\epsilon \in \mathcal{E}} \mathcal{U}_{s}(\nu \lfloor (\mathbb{R}^{p} \setminus B(y_{\infty}, \delta)), y_{\infty, \epsilon}) \geq \mathcal{U}_{s}(\nu \lfloor (\mathbb{R}^{p}, \backslash B(y_{\infty}, \delta)), y_{\infty}). \tag{11}$$

After replacing the sequences \mathcal{E} , \mathcal{N} by the subsequences \mathcal{E}' , \mathcal{N}' along which the liminf is achieved, we now take the limit over $\delta \to 0$. Again by the principle of descent, from (11) we get

$$\lim_{\delta \to 0} \lim_{\epsilon \in \mathcal{E}'} \lim_{N \in \mathcal{N}'} \mathcal{U}_{s}(\nu_{N} \lfloor (\mathbb{R}^{p} \setminus B(y_{\infty}, \delta)), y_{N, \epsilon})$$

$$\stackrel{(11)}{\geq} \lim_{\delta \to 0} \mathcal{U}_{s}(\nu \lfloor (\mathbb{R}^{p} \setminus B(y_{\infty}, \delta)), y_{\infty}) = \mathcal{U}_{s}(\nu, y_{\infty}) \geq P_{s}(A, \nu).$$
(12)

To justify the above equality, note that the potentials $U_s(\nu \mid (\mathbb{R}^p \setminus B(y_\infty, \delta)), y)$ are decreasing in δ and that ν has no atoms.

From (9), (10) and (12) and since for any $\delta > 0$ there holds $\mathcal{U}_s(\nu_{N,\epsilon}, y) \ge \mathcal{U}_s(\nu_{N,\epsilon,\delta}, y)$ at all $y \in \mathbb{R}^d$, we get that

$$\lim_{\epsilon \in \mathcal{E}'} \lim_{N \in \mathcal{N}'} P_{s}(A, \nu_{N, \epsilon}) = \lim_{\epsilon \in \mathcal{E}'} \lim_{N \in \mathcal{N}'} \mathcal{U}_{s}(\nu_{N, \epsilon}, y_{N, \epsilon})$$

$$\geq \lim_{\delta \to 0} \lim_{\epsilon \in \mathcal{E}'} \lim_{N \in \mathcal{N}'} \mathcal{U}_{s}(\nu_{N, \epsilon, \delta}, y_{N, \epsilon}) \geq P_{s}(A, \nu). \tag{13}$$

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We now add $\#(\omega_N^A \setminus A_{\epsilon})$ points from A to the $\omega_{N,\epsilon}^A$, obtaining new configurations with N points. This has the only effect of increasing the generated potential. Since we can apply all the above reasoning to any initial subsequence $\mathcal{N}^0 \subset \mathbb{N}$, we get

$$\liminf_{N \to \infty} \frac{\mathcal{P}_s(A, N)}{N} \ge P_s(A, \nu).$$
(14)

On the other hand, the unconstrained optimum polarization is always larger than the constrained optimum polarization, and by using (6) we have

$$P_s(A, \nu) = \lim_{N \to \infty} \frac{\mathcal{P}_s^*(A, N)}{N} \ge \limsup_{N \to \infty} \frac{P_s(A, N)}{N}.$$
 (15)

From (14) and (15) it follows that

$$\lim_{N \to \infty} \frac{\mathcal{P}_s^*(A, N)}{N} = \lim_{N \to \infty} \frac{\mathcal{P}_s(A, N)}{N},\tag{16}$$

as desired.

We now address Theorem 2 and its proof. The authors thank Alex Vlasiuk for pointing out that the derivation of equation (5.23) from equation (5.22) of [3] in the proof of [3, Thm. 1.12] did not take into account the measure of the boundary of A. We provide in Proposition 4 below, a substitute for this derivation valid in the case s > p, whereas for s = p we add to [3, Thm. 1.12] the additional hypothesis $\mathcal{L}^p(\partial A) = 0$.

Proof of Theorem 2. The proof of Theorem 2 follows exactly like the one of [3, Thm. 1.12], except for the following changes:

- For the case s = p, with the further hypothesis $\mathcal{L}^p(\partial A) = 0$ in Theorem 2, we can take the sets $G_i := A \cap B_i$ in the paragraph following (5.21) of [3], and the proof holds verbatim.
- For the case s > p, Proposition 4 below, applied to the sets $A \cap B_i$ and B_i from eq. (5.22) of [3] allows to replace eq. (5.23) therein. Observing that B is convex and thus $h_{s,p}^*(B) = h_{s,p}(B)$ by [3, Prop. 1.7], and that by [1] we have $h_{s,p}(B)^{-p/s} = (\sigma_{s,p})^{-p/s}\mathcal{L}_p(B)$, it follows that $\bar{h}_{s,p}^*(B) = (\sigma_{s,p})^{-p/s}\mathcal{L}_p(B)$. The factor $1 2\epsilon$ in eq. (5.23), (5.24) of [3] should be replaced by $(1 c_{s,p}\epsilon^{\frac{1}{p+1}})^{p/s}$ with notations as in Proposition 4 below. Noting that both these quantities tend to 1 as $\epsilon \to 0$, the proof of Theorem 1.12 of [3] follows with no further modifications.

The new result needed for the case s > p in Theorem 2 is the following:

Proposition 4 For $s > p \ge 1$, there exists a constant $c_{s,p} > 0$ with the following properties. Let $\epsilon \in (0, 1/2^{p+1})$ and $B \subset \mathbb{R}^p$ be a ball and $A \subset B$ a closed set such that $\mathcal{L}^p(B \setminus A) < \epsilon \mathcal{L}^p(B)$. Then there holds

$$\overline{h}_{s,p}^*(A) \le (1 - c_{s,p} \epsilon^{\frac{1}{p+1}})^{-1} \overline{h}_{s,p}^*(B). \tag{17}$$

In what follows we set $A, B \subset \mathbb{R}^p$ to be as in the statement of Proposition 4. Furthermore, let $\mathcal{P}_s^*(A, N)$ be the optimum N-point polarization of set A, and let ω_N^A be an N-point configuration be as in (5). The proof of Proposition 4 is based on two lemmas.

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Lemma 5 Let $0 < \delta < \min\{1, s/2\}$ and let N be a positive integer. If $y \in \mathbb{R}^p$ is such that

$$\operatorname{dist}(y,A) < \frac{\delta}{s} \left(\mathcal{P}_s^*(A,N) \right)^{-\frac{1}{s}} \tag{18}$$

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then

$$\mathcal{U}_s(\omega_N^A, y) \ge (1 - 2\delta)\mathcal{P}_s^*(A, N). \tag{19}$$

Proof First note that if $\operatorname{dist}(y, \omega_N^A) < (\mathcal{P}_s^*(A, N))^{-\frac{1}{s}}$ then

$$\mathcal{U}_s(\omega_N^A, y) \geq \mathcal{P}_s^*(A, N),$$

and thus (19) holds a fortiori. Therefore from now on we consider points $y \in \mathbb{R}^p$ such that (18) and $\operatorname{dist}(y, \omega_N^A) \geq (\mathcal{P}_s^*(A, N))^{-\frac{1}{s}}$ hold, and our goal is to prove (19) for such y. Note that as a consequence of this assumption and of (18), we also get

$$\operatorname{dist}(y, A) \le \frac{\delta}{s} \operatorname{dist}(y, \omega_N^A). \tag{20}$$

Let $y_1 \in A$ be such that $|y_1 - y| = \text{dist}(y, A)$ and let

$$y_2 \in \operatorname{argmax}_{y' \in [y, y_1]} \mathcal{U}_s(\omega_N^A, y').$$
 (21)

Note that if $y' \in [y, y_1]$ and x_j is the point from ω_N^A closest to y', we have, using (20),

$$\operatorname{dist}(y', \omega_N^A) = |y' - x_j| \ge |y - x_j| - |y' - y| \ge \operatorname{dist}(y, \omega_N^A) - \operatorname{dist}(y, A) \ge \left(1 - \frac{\delta}{s}\right) \operatorname{dist}(y, \omega_N^A). \tag{22}$$

We now claim that the following chain of inequalities holds:

$$\mathcal{U}_{s}(\omega_{N}^{A}, y) \geq \mathcal{U}_{s}(\omega_{N}^{A}, y_{2}) - \int_{0}^{1} \left| \nabla \mathcal{U}_{s}(\omega_{N}^{A}, y + t(y_{2} - y)) \cdot (y_{2} - y) \right| dt \tag{23}$$

$$\geq \mathcal{U}_s(\omega_N^A, y_2) - s|y_2 - y| \left[\min_{y' \in [y, y_2]} \operatorname{dist}(y', \omega_N^A) \right]^{-1} \max_{y' \in [y, y_2]} \mathcal{U}_s(\omega_N^A, y') \quad (24)$$

$$\geq (1 - 2\delta) \mathcal{U}_s(\omega_N^A, y_2) \geq (1 - 2\delta) \mathcal{P}_s^*(A, N). \tag{25}$$

We now prove the above. The bound (23) follows by Taylor expansion. Inequality (24) follows by noting that whenever $x \neq y'$ we have $|\nabla_y|x - y'|^{-s}| = s|x - y'|^{-s-1}$, therefore for $y' \in [y, y_2]$ we have

$$\left|\nabla \mathcal{U}_s(\omega_N^A, y')\right| \leq s \sum_{x \in \omega_N^A} |x - y'|^{-s - 1} \leq s \left[\min_{y' \in [y, y_2]} \operatorname{dist}(y', \omega_N^A) \right]^{-1} \mathcal{U}_s(\omega_N^A, y').$$

For the first inequality in (25), note that from our hypotheses on y we get $|y_2 - y| \le |y_1 - y| = \operatorname{dist}(y, A) \le \frac{\delta}{s} (\mathcal{P}_s^*(A, N))^{-\frac{1}{s}}$ and from (22) we get $\min_{y' \in [y, y_2]} \operatorname{dist}(y', \omega_N^A) \ge (1 - \frac{\delta}{s}) (\mathcal{P}_s^*(A, N))^{-\frac{1}{s}}$. This gives, using the hypothesis $\delta < s/2$,

$$s|y_2 - y| \left[\min_{y' \in [y, y_2]} \operatorname{dist}(y', \omega_N^A) \right]^{-1} \le \frac{\delta s}{s - \delta} < 2\delta.$$

To get the first inequality in (25) it now suffices to recall definition (21) of y_2 . The second inequality in (25) follows by the fact that $y_1 \in A$ and the definition of y_2 , and of ω_N^A :

$$\mathcal{P}_{s}^{*}(A, N) = \min_{x \in A} \mathcal{U}_{s}(\omega_{N}^{A}, x) \leq \mathcal{U}_{s}(\omega_{N}^{A}, y_{1}) \leq \mathcal{U}_{s}(\omega_{N}^{A}, y_{2})$$

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Recall the notation, for the r-neighborhood of a closed set $K \subset \mathbb{R}^d$, for r > 0:

$$(K)_r := \{ x \in \mathbb{R}^d : \operatorname{dist}(x, K) < r \}.$$

Lemma 6 There exists a constant $c_p > 0$ with the following properties. Let $B \subset \mathbb{R}^p$ be a ball and let $A \subset B$ be a closed set, such that for some $\epsilon \in (0,1)$ there holds $\mathcal{L}^p(B \setminus A) < \epsilon \mathcal{L}^p(B)$. Then for each $r \in (0,\epsilon \text{ diam}(B))$ we can cover all of $B \setminus (A)_r$ by at most $c_p \frac{\epsilon \mathcal{L}^p(B)}{r^p}$ balls of radius r.

Proof Let $r' := r/(2\sqrt{p})$ and we show that we may take as the set of ball centers the following:

$$W := \{r'k : k \in \mathbb{Z}^p, (r'k + [-r', r']^p) \cap (B \setminus (A)_r) \neq \emptyset\}.$$

We note that the balls with centers in W and radius $\frac{r'}{2}$ are disjoint and contained in the r-neighborhood of $B \setminus (A)_r$. Furthermore, we have the inclusion

$$(B \setminus (A)_r)_r \subset (B)_r \setminus A = ((B)_r \setminus B) \cup (B \setminus A),$$

from which it follows that, denoting by #W the cardinality of W,

$$\#W \cdot \mathcal{L}^p(B_{\underline{r'}}) \leq \mathcal{L}^p((B)_r \setminus B) + \mathcal{L}^p(B \setminus A) \leq \left(C_p \frac{r}{\operatorname{diam}(B)} + \epsilon\right) \mathcal{L}^p(B) \leq (C_p + 1) \epsilon \mathcal{L}^p(B).$$

This implies that for $c_p := 2^p (C_p + 1)/(p^{\frac{p}{2}}\beta_p)$, in which β_p is the volume of the unit ball in \mathbb{R}^p , there holds

$$\#W \le c_p \epsilon \frac{\mathcal{L}^p(B)}{r^p}.$$

It remains to show that radius-r balls with centers in W cover $B \setminus (A)_r$. Indeed, note that if the cube $r'k + [-r', r']^p$ with $k \in \mathbb{Z}^p$ meets $B \setminus (A)_r$, then it is contained in the ball $B(r'k, \sqrt{p} \ r') \subset B(r'k, r)$, and thus

$$B \setminus (A)_r \subset \bigcup_{r'k \in W} (r'k + [-r', r']^p) \subset \bigcup_{r'k \in W} B(r'k, r),$$

as desired.

Proof of Proposition 4: Observe that by the same proof as in [2, Thm. 2.4], which applies also without the restriction that the optimum polarization points belong to A, with \mathcal{L}^p used as the measure μ in the proof, there exists a constant $C_{s,p} > 0$ independent of N, A, such that $\mathcal{P}_s^*(A,N) \leq C_{s,p} N^{s/p} / (\mathcal{L}^p(A))^{s/p}$. Now, applying Lemma 5 with $\delta := \epsilon^{\frac{1}{p+1}}$ we find that for the optimum configuration ω_N^A like in (5), with

$$r_{N} := \frac{\epsilon^{\frac{1}{p+1}}}{s} \left(\mathcal{P}_{s}^{*}(A, N) \right)^{-\frac{1}{s}} \ge \frac{\epsilon^{\frac{1}{p+1}} \left(\mathcal{L}^{p}(A) \right)^{\frac{1}{p}}}{s(C_{s, p})^{\frac{1}{s}}} N^{-\frac{1}{p}}, \tag{26}$$

we have

$$\forall y \in (A)_{r_N}, \quad \mathcal{U}_s(\omega_N^A, y) \ge \left(1 - 2\epsilon^{\frac{1}{p+1}}\right) \mathcal{P}_s^*(A, N).$$
 (27)

Next, for N large enough so that $r_N < \epsilon \operatorname{diam}(B)$ we apply Lemma 6 with r_N playing the role of r, and we find a set of centers W such that, using also (26) in the second below

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inequality:

$$B\setminus (A)_{r_N} \subset \bigcup_{x\in W} B(x,r_N), \quad \#W \leq c_p \, \epsilon \frac{\mathcal{L}^p(B)}{r_N^p} \leq \widetilde{C}_{s,p} \, \epsilon^{\frac{1}{p+1}} \, \frac{\mathcal{L}^p(B)}{\mathcal{L}^p(A)} \, N \leq \widetilde{C}_{s,p} \, \frac{\epsilon^{\frac{1}{p+1}}}{1-\epsilon} N.$$
(28)

By considering the new configuration $\omega_{\widetilde{N}} := \omega_N^A \cup W$ whose cardinality is denoted \widetilde{N} , we

$$\widetilde{N} \in \left[N, N\left(1 + \widetilde{C}_{s,p} \frac{\epsilon^{\frac{1}{p+1}}}{1 - \epsilon}\right)\right], \quad \mathcal{U}_{s}(\omega_{\widetilde{N}}, y) \ge \begin{cases} \left(1 - 2\epsilon^{\frac{1}{p+1}}\right) \mathcal{P}_{s}^{*}(A, N) & \text{for } y \in (A)_{r_{N}}, \\ \mathcal{P}_{s}^{*}(A, N) & \text{for } y \in B \setminus (A)_{r_{N}}. \end{cases}$$
(29)

The bound in the first line of (29) follows from (27). For the second bound in (29), note that since $\epsilon < 1 \le p < s$, using the first part of (28) the following holds for $x \in W$, $y \in$ $B(x, r_N)$:

$$\mathcal{U}_s(\omega_{\widetilde{N}}, y) \ge |x - y|^{-s} \ge r_N^{-s} = \left(\frac{\epsilon^{\frac{1}{p+1}}}{s}\right)^{-s} \mathcal{P}_s^*(A, N) \ge \mathcal{P}_s^*(A, N).$$

We then find that, using also the fact that $\epsilon \in (0, 1/2^{p+1})$

$$\frac{\mathcal{P}_{s}^{*}(B,\widetilde{N})}{\widetilde{N}^{s/p}} \geq \frac{1 - 2\epsilon^{\frac{1}{p+1}}}{\left(1 + \widetilde{C}_{s,p} \frac{\epsilon^{\frac{1}{p+1}}}{1 - \epsilon}\right)^{s/p}} \frac{\mathcal{P}_{s}^{*}(A,N)}{N^{s/p}} \geq (1 - 2\epsilon^{\frac{1}{p+1}}) \left(1 - \frac{s}{p} \widetilde{C}_{s,p} \frac{\epsilon^{\frac{1}{p+1}}}{1 - \epsilon}\right) \frac{\mathcal{P}_{s}^{*}(A,N)}{N^{s/p}},$$

(30)from which the bound (17) follows, with $c_{s,p} = 2 + \frac{2^{p+1}}{2^{p+1}-1} \frac{s}{p} \widetilde{C}_{s,p}$.

Now the proof of Theorem 2 is complete.

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References

- 1. Borodachov, S.V., Hardin, D.P., Reznikov, A., Saff, E.B.: Optimal discrete measures for Riesz potentials. Trans. Amer. Math. Soc. 370(10), 6973-6993 (2018)
- 2. Erdélyi, T., Saff, E.B.: Riesz polarization inequalities in higher dimensions. J. Approx. Theory 171, 128-147 (2013)
- 3. Hardin, D.P., Petrache, M., Saff, E.B.: Asymptotics for the unconstrained polarization (Chebyshev) problem. Potential Anal, https://doi.org/10.1007/s11118-020-09875-z
- 4. Ohtsuka, M.: On various definitions of capacity and related notions. Nagoya Math. J. 30, 121–127 (1967)
- 5. Reznikov, A., Saff, E.B., Vlasiuk, O.V.: A minimum principle for potentials with application to Chebyshev constants. Potential Anal. 47(2), 235–244 (2017)
- Simanek, B.: Asymptotically optimal configurations for Chebyshev constants with an integrable kernel. New York J. Math. 22, 667–675 (2016)

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56, 21–64 (2022). https://doi.org/10.1007/s11118-02 0-09875-z

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